

Comment on “Universal Fluctuations in Correlated Systems”

In [1], it is suggested that the probability distribution functions (PDF's) of fluctuating observables for critical systems in different universality classes exhibit the same form. This Comment concerns the PDF's for equilibrium systems.

In [2], it is pointed out that the standard scaling form is a sufficient condition for the data collapse. The suggestion in [1] seems to contradict the idea of standard universality. To clarify it, we have performed calculations for different critical systems. First, the temperature is fixed at the bulk critical value T_c . In Fig. 1, the normalized PDF $P(m)$ for the Ising and XY models are displayed. Here $m = (M - \langle M \rangle) / \sigma$, where M is the magnetization, $\langle M \rangle$ is its mean, and σ is the standard deviation. Data collapse for different lattice sizes L 's is observed. However, the differences between the curves for different models are clearly not minor perturbations. The PDF for the Potts model is even not symmetric in M and can hardly be put in the figure.

The critical temperature can be modified in a finite system. Let us take a size-dependent coupling $K(L) \sim 1/T(L)$ such that $\tau \equiv [K(L) - K_c] / K_c = s/L^{1/\nu}$, where ν is the static exponent and s is a constant. Following similar scaling analysis as in [2] (see also [3]), it leads to $P(m, \tau, L) = P(m, L^{1/\nu} \tau) = P(m, s)$. At the critical regime, data for different L 's at $T(L)$'s also collapse onto a single curve, but $P(m, s)$ changes continuously with s . If s is small, in other words, if the spatial correlation length $l(L)$ is much larger than the lattice size L , $T(L)$ can be considered approximately as a size-dependent critical temperature. Choosing $s = 2.90$, $P(m, s)$ for the 2D Ising model falls onto that for the 2D XY model at $T = 0.89$. This is shown in Fig. 1. However, $T(L)$ with such a large value of s should not be defined as a size-dependent critical temperature, since in the infinite limit of L , the behavior of the system [not only $P(m)$] remains very different from that at the bulk T_c . To confirm this, we have calculated the spatial correlation function and found that at $s = 2.90$ the correlation length $l(L)$ is much smaller than the lattice size L .

We do not think the *shape* of $P(m)$ for the XY model is a characteristic property at the critical point. The observation is that for systems with a second order transition, the tail of the PDF for negative m at the T_c reaches a nonzero value at $M = 0$ and is roughly power-law-like (before cut at $M = 0$). This reasonably indicates that the system can transit from the positive sector of M to the negative one. Below T_c , symmetry breaking occurs. The (infinite) system cannot transit from one sector to another. The exponential-like tail at a large but not too large s should be a signal of symmetry breaking. (But when s is infinite, the tail crosses over to Gaussian.) For the XY model, the

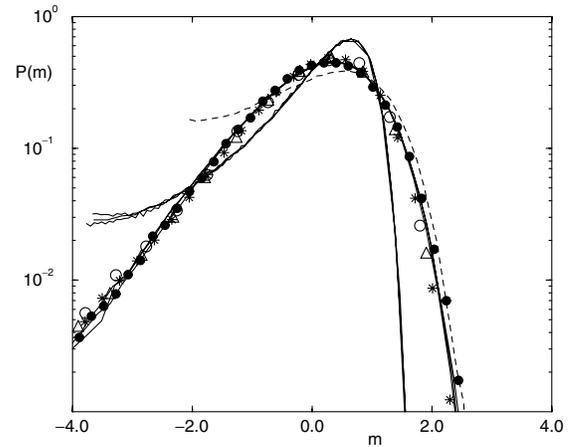


FIG. 1. The three solid lines with a higher peak are for the 2D Ising model at T_c with $L = 32, 64,$ and 128 . The three solid lines with a lower peak are for the 2D XY model at $T = 0.89$ with $L = 16, 32,$ and 64 . The dashed curve is for the 3D Ising model at T_c with $L = 32$. The circles, triangles, and stars are for the 2D Ising model at $s = 2.90$ with $L = 32, 64,$ and 128 . The filled circles are for the 3D Ising model at $s = 2.21$ with $L = 32$.

fluctuations are mainly rotational. The exponential tail of $P(m)$ is only an indication for the energy barrier in small M regime. The conjecture is that the exponential-like tail of the PDF induced typically by energy barriers may be similar to the exponential decay of the correlation functions with a finite correlation length. It can be rather generic and independent of whether the system is with or without a first order, second order, or Kosterlitz-Thouless phase transition. The results for the 1D and 3D XY models in [3] support this statement. A critical point is only a sufficient condition for data collapse for infinite systems.

The complete form of the PDF is, in general, not independent of universality classes. However, the large s regime may be somewhat special and it needs more investigation. For example, as shown in Fig. 1, $P(m, s)$ for the 3D Ising model at $s = 2.21$ fits also to the curve of the XY model at $T = 0.89$.

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